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Phase dynamics in the real Ginzburg-Landau equation

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Abstract. Spatially periodic equilibria $A(X,T) = \sqrt{1-q^2}e^{iqX+i\phi_0}$ are the locally preferred planform for the Ginzburg-Landau equation $\partial_T A = \partial_X^2 A + A - A|A|^2$. To describe the global spatial behavior, an evolution equation for the local wave number q can be derived formally. The local wave number q satisfies approximately a so called phase diffusion equation $\partial_\tau q = \partial_\xi^2 h(q)$. It is the purpose of this paper to explain the extent to which the phase diffusion equation is valid by proving estimates for this formal approximation.

1. Introduction

The (normalized) real Ginzburg-Landau equation

$$\partial_T A = \partial_X^2 A + A - A|A|^2$$

with $X \in \mathbb{R}$, $T \geq 0$, and $A(X,T) \in \mathbb{C}$, is a universal amplitude equation which is derived by multiple scaling analysis in order to describe bifurcating solutions in pattern forming systems close to the threshold of the first instability. The amplitude A describes slow modulations in space and time of the underlying bifurcating spatially periodic pattern. Examples of such pattern forming systems are reaction-diffusion systems, systems in nonlinear optics, or hydrodynamic stability problems, for example Rayleigh-Bénard convection or the Taylor-Couette problem. A mathematical theory of this reduction to the real (and complex) Ginzburg-Landau equation has been developed by several authors (cf. [CE90, vH91, Sch94, Me98, Me99]). It is nowadays a well established mathematical tool which can be used to obtain new mathematical results (cf. [Sch99]), such as global existence results and uppersemicontinuity of attractors.

The real Ginzburg-Landau equation possesses a family of stationary solutions

$$A(X,T) = \sqrt{1-q^2}e^{iqX+i\phi_0} = A_{\mathrm{per}}[q,\phi_0](X)$$

with $q \in (-1,1)$, $\phi_0 \in \mathbb{R}$. For $q^2 \leq 1/3$ the equilibria are spectrally stable, whereas for $q^2 > 1/3$ they are Eckhaus or sideband unstable [Eck65]. Locally in space, these

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solutions are the preferred planform. In order to describe the global behavior in space, an evolution equation for the local wave number q can be derived. Allowing q to vary slowly in time and space, we write

$$A_{\text{per}}[\psi(\delta X, \delta^2 T), \phi_0](X) := \sqrt{1 - (\psi(\delta X, \delta^2 T))^2} \exp\Big(i \int_0^X \psi(\delta X', \delta^2 T) dX' + i\phi_0\Big)$$

with $0 < \delta \ll 1$ a small perturbation parameter, where ψ satisfies the phase diffusion equation

$$\partial_\tau \psi = \partial_\xi^2 h(\psi) \tag{1.1}$$

with $\tau = \delta^2 T$, $\xi = \delta X$, and $h : \mathbb{R} \to \mathbb{R}$ a smooth function. It is the purpose of this paper to explain to which extent this formal approximation is valid by proving estimates between the formal approximation $A_{\text{per}}[\psi(\delta X, \delta^2 T), \phi_0](X)$ and exact solutions $A = A(X,T)$ of the real Ginzburg-Landau equation.

It turns out that we cannot expect validity uniformly for all $X \in \mathbb{R}$, but validity only uniformly for all $X \in I_\delta$ with I_δ an interval of length $O(\delta^{-r})$ for all $r < 2$.

Other amplitude equations for the evolution of the local wave number of stable and unstable planforms close to the Eckhaus-stable threshold in the Ginzburg-Landau equation have been considered in [Be88, vH95].

It is not obvious *a priori* that an approximation result for the phase diffusion equation (1.1) holds. There are a number of counterexamples of amplitude equations which are derived formally in a correct way, but do not describe the dynamics in the original system in a correct way. See for instance [Sch95].

The difficulty in justifying the phase-diffusion equation for the Ginzburg-Landau equation is the time scale $O(1/\delta^2)$. Since the solutions in consideration are of order $O(1)$ a simple application of Gronwall's inequality would only give a time scale $O(1)$. In order to gain estimates on the long time scale $O(1/\delta^2)$ the proof is built similarly to the derivation of the phase diffusion equation. In particular, in polar coordinates the equations for the error turn out to be quasilinear.

Our approximation result allows us to find the dynamics of the phase diffusion equation in the real Ginzburg-Landau equation. Moreover, the Ginzburg-Landau equation approximates more complicated pattern forming systems like the Taylor-Couette problem, close to the first instability, and so we can find the dynamics of the phase diffusion equation in these more complicated systems.

The dynamics of nonlinear diffusion equations like the phase diffusion equation is well studied. There are a number of self similar solutions [BKL94]. Moreover, in contrast to the original systems, the maximum principle holds.

Away from the threshold of the first instability, phase-diffusion equations for the evolution of the local wave number can be derived in order to describe spatial and temporal modulations of the fully developed spatially periodic pattern. Formally, they have been derived and discussed by several authors in a number of papers (cf. [CN84, CH93, EINP00, MP79, Ma90, NPL93, PN94, KP96]). There are also mathematically rigorous stability and diffusive mixing results which show that at least some of the dynamics predicted by the phase-diffusion equations can be found in the Ginzburg-Landau equation (cf. [CEE92, BK92, GM98]) or in the associated pattern forming systems (cf. [Sch98]). It is the purpose of further research to justify the phase-diffusion equations themselves away from the threshold of the first instability.

The plan of the paper is as follows. In Section 2 we derive the phase-diffusion equation by introducing polar coordinates $A = re^{i\phi}$ and introducing $\psi = \partial_X \phi$. In Section 3 we prove estimates which hold uniformly in space for the variables (r, ψ). In Section 4 we go back to the original A-variable which leads to the result that estimates which hold uniformly in space cannot be expected for the approximation of A. We also explain that the uniformity of the approximation can be improved by solving the equations for the higher order corrections. See Remark 4.2. We note that the alternative approach of [Me98, Me99], discussed briefly in Remark 3.5, shows that the derivation of the phase-diffusion equation can be made exact for a certain class of solutions if derivative terms of all orders are included (so that equation (1.1) becomes a pseudodifferential equation).

Notation. Throughout this paper many different constants are denoted with the same symbol C. We denote the space of n-times weakly differentiable local uniformly Sobolev functions with $H^n_{l,u}$. This Banach space is equipped with the norm

$$\|u\|_{H^n_{l,u}} = \sup_{x\in\mathbb{R}} \|u(\cdot)\|_{H^n(x,x+1)}.$$

For details we refer to [Sch99].

2. Derivation of the phase-diffusion equation for the real Ginzburg-Landau equation

As already said, the purpose of this paper is to justify a phase-diffusion equation describing the evolution of the wave number q of the spatially locally preferred planform for the real Ginzburg-Landau equation

$$\partial_T A = \partial_X^2 A + A - A|A|^2 \tag{2.1}$$

with $X \in \mathbb{R}$, $T \geq 0$, and $A(X,T) \in \mathbb{C}$. This equation possesses a family of stationary solutions

$$A = A_{\text{per}}[q, \phi_0](X) = \sqrt{1-q^2} e^{iqX + i\phi_0} \tag{2.2}$$

with $\phi_0, q \in \mathbb{R}$ as spatially locally preferred planform. To derive the phase diffusion equation we introduce polar coordinates

$$A(X,T) = r(X,T)e^{i\phi(X,T)}$$

in the real Ginzburg-Landau equation and obtain

$$\left\{ \begin{array}{rcl} \partial_T r &=& \partial_X^2 r + r - (\partial_X \phi)^2 r - r^3, \\ \partial_T \phi &=& \partial_X^2 \phi + \dfrac{2(\partial_X r)(\partial_X \phi)}{r}. \end{array} \right. \tag{2.3}$$

Besides the interpretation in the introduction, another interpretation will be that the phase-diffusion equation is derived to describe the dynamics close to a stationary pattern $A_{\text{per}}[q, \phi_0]$. We choose $(r, \phi) = (1, 0)$ as the starting stationary pattern. For the general case see Remark 2.1 below. We introduce the deviations (s,ϕ) by setting $r=1+s$ to obtain

$$\begin{aligned}\partial_T s &= \partial_X^2 s - 2s - (\partial_X\phi)^2 - (\partial_X\phi)^2 s - 3s^2 - s^3,\\ \partial_T\phi &= \partial_X^2\phi + \frac{2(\partial_X s)(\partial_X\phi)}{1+s}.\end{aligned}$$

We can replace the equation for the phase ϕ by an equation for the local wave number $\psi=\partial_X\phi$ to obtain

$$\left\{\begin{aligned}\partial_T s &= \partial_X^2 s - 2s - \psi^2 - \psi^2 s - 3s^2 - s^3,\\ \partial_T\psi &= \partial_X^2\psi + 2\partial_X\left(\frac{(\partial_X s)\psi}{1+s}\right).\end{aligned}\right. \tag{2.4}$$

The linearized system

$$\begin{aligned}\partial_T s &= \partial_X^2 s - 2s,\\ \partial_T\psi &= \partial_X^2\psi,\end{aligned}$$

possesses solutions $(s,\psi)=(s_k,\psi_k)e^{ikx+\mu_j(k)t}$ with $\mu_1(k)=-k^2$, and $\mu_2(k)=-k^2-2$. Hence, the component s which corresponds to μ_2 is exponentially damped and so we expect s to be slaved by ψ which will behave diffusively. To derive the phase-diffusion equation we make a long wave ansatz

$$\psi(X,T)=\check\psi(\delta X,\delta^2T) \qquad \text{and} \qquad s(X,T)=\check s(\delta X,\delta^2 T)$$

with $0\le\delta\ll 1$ a small perturbation parameter. This yields

$$\left\{\begin{aligned}\delta^2\partial_\tau\check s &= \delta^2\partial_\xi^2\check s - 2\check s - \check\psi^2 - \check\psi^2\check s - 3\check s^2 - \check s^3,\\ \partial_\tau\check\psi &= \partial_\xi^2\check\psi + 2\partial_\xi\left(\frac{\check\psi\partial_\xi\check s}{1+\check s}\right),\end{aligned}\right. \tag{2.5}$$

with $\tau=\delta^2T$ and $\xi=\delta X$. Neglecting terms of order $O(\delta^2)$ for small $\check\psi$ by the implicit function theorem, the first equation

$$0=-2\check s-\check\psi^2-\check\psi^2\check s-3\check s^2-\check s^3 \tag{2.6}$$

possesses a unique solution $\check s=s^*(\check\psi)$ with $s^*(\check\psi)=-\check\psi^2/2+O(\check\psi^3)$. (Due to the formula for the equilibria, we have in fact that $1+s^*=\sqrt{1-\check\psi^2}$.) Inserting this into the equation for $\check\psi$ leads finally to the phase diffusion equation

$$\partial_\tau\check\psi=\partial_\xi^2\check\psi+2\partial_\xi\left(\frac{\check\psi\partial_\xi(s^*(\check\psi))}{1+s^*(\check\psi)}\right)=\partial_\xi^2 h(\check\psi) \tag{2.7}$$

where $h:\mathbb{R}\to\mathbb{R}$ is smooth. Indeed, h is the solution to the equation

$$h'(q)=1+\frac{2q}{1+s^*(q)}s^{*\prime}(q)=\frac{1-3q^2}{1-q^2}.$$

Remark 2.1. Suppose that instead we start with a general choice of basic equilibrium $r=r_q$, $\phi=qX+\phi_0$, $q\in(-1,1)$. Then the corresponding phase diffusion equation is given by

$$\partial_\tau\check\psi=\partial_\xi^2\tilde h(\check\psi,q) \tag{2.8}$$

where $\tilde{h}(\check{\psi}, q) = h(\check{\psi} + q)$.

To verify this, note that the deviations $\check{\psi}_q$ for (2.8) are related to the deviations $\check{\psi}_0$ for (2.7) by $\check{\psi}_q + q = \check{\psi}_0$. Hence $\partial_\tau \check{\psi}_q = \partial_\tau \check{\psi}_0 = h(\check{\psi}_0) = h(\check{\psi} + q) = \tilde{h}(\check{\psi}, q)$ as required.

Remark 2.2. For each q, let $k \mapsto \mu_{1,2}(k, q)$ denote the smooth curves of eigenvalues corresponding to the Fourier wave numbers k for the linearization of (2.3) around $(r, \phi) = (\sqrt{1 - q^2}, qX)$. In particular, let $\mu_1(k, q)$ denote the critical curve for which $\mu_1(0, q) = 0$. Then we claim that the phase-diffusion equation (2.7) must give at lowest order a linear diffusion equation with diffusion coefficient

$$h'(q) = -\partial_k^2 \mu_1(0, q)/2.$$

This is easily verified by direct calculation, since we have already seen that $h'(q) = (1 - 3q^2)/(1 - q^2)$ while $\mu_1(k, q) = -k^2 - (1 - q^2) + \sqrt{(1 - q^2)^2 + 4q^2k^2}$. A more conceptual proof is as follows.

First note that the linearization $L(q)$ of the right-hand-side $\partial_X^2 \tilde{h}(\psi, q)$ of (2.8) at $(0, q)$ must also have the eigenvalues $\mu_1(k, q)$ – after taking into account the fact that higher order derivatives have been neglected in the derivation of (2.8). But

$$\partial_X^2 \tilde{h}(\psi, q) = \partial_X^2 h(q + \psi) = h'(q)\partial_X^2 \psi + O(|\psi|^2)$$

and so $L(q) = h'(q)\partial_X^2$. Equating $L(q)e^{ikX} = \mu_1(k, q)e^{ikx}$ modulo higher order derivatives yields $h'(q) = -\partial_k^2 \mu_1(0, q)/2$ as required.

Note that $\partial_k^2 \mu_1(0, q) > 0$ if and only if $q^2 > 1/3$ which is the famous Eckhaus instability criterion [Eck65].

Remark 2.3. If the solutions of the phase-diffusion equation (2.7) start in between the Eckhaus stable range, i.e. having values $|\check{\psi}|^2 < 1/3$, they stay in this Eckhaus stable range for all times due to the maximum principle. For sufficiently smooth initial conditions satisfying $\sup_{X \in \mathbb{R}} |\check{\psi}(X)|^2 < 1/3$ there is the local and global existence of solutions (cf. [Kn68, LSU68, Va92]).

Remark 2.4. It is common in the literature to consider generalized versions of the real Ginzburg-Landau equation with more complicated nonlinearities. For example, terms such as $i\partial_X A|A|^2$ and $|A|^4 A$ are often included. More generally, terms of the form $A^{b_1}\bar{A}^{b_2}(i\partial_X A)^{b_3}(i\partial_X \bar{A})^{b_4}$ are permitted provided $b_1 - b_2 + b_3 - b_4 = 1$. In this case, writing $A = re^{i\phi}$, $r = 1 + s$, $\psi = \partial_X \phi$, leads to a system of the following form in place of (2.4):

$$\left\{ \begin{array}{rcl} \partial_T s &=& \partial_X^2 s + f(s, (\partial_X s)^2, \psi), \\ \partial_T \psi &=& \partial_X^2 \psi + \partial_X \Big((\partial_X s) g(s, (\partial_X s)^2, \psi) \Big) \end{array} \right. \tag{2.9}$$

where f and g are smooth functions of $s, \psi, (\partial_X s)^2$.

This structure can be verified as follows. The symmetries $(A \mapsto \bar{A}, X \mapsto -X)$ and $A \mapsto e^{ic}A$ of the real Ginzburg-Landau equation lead to the symmetries $(\phi \mapsto -\phi, X \mapsto -X)$ and $\phi \mapsto \phi + c$ for the (s, ϕ) equations. Hence, we obtain $\partial_T s =$

$\partial_X^2 s + f(s, (\partial_X s)^2, \partial_X \phi)$, $\partial_T \phi = \partial_X^2 \phi + (\partial_X s) g(s, (\partial_X s)^2, \partial_X \phi)$. Writing $\psi = \partial_X \phi$ yields (2.9).

Now we can write $\tau = \delta^2 T$ and $\xi = \delta X$ and neglecting small terms we reduce as before to the phase-diffusion equation

$$\partial_\tau \check{\psi} = \partial_\xi^2 \check{\psi} + \partial_\xi \Big((\partial_\xi s^*(\check{\psi})) g(s^*(\check{\psi}), 0, \check{\psi}) \Big) = \partial_\xi^2 h(\check{\psi}),$$

where h satisfies $h'(q) = 1 + {s^*}'(q) g(s^*(q), 0, q)$. Again, $h'(q) = -\partial_k^2 \mu_1(0,q)/2$ where $k \mapsto \mu_1(k,q)$ is the smooth critical curve of eigenvalues corresponding to Fourier wave numbers k. The estimates of this paper can be proved in a similar manner for this more general system, too.

3. The approximation theorem for the $(\check{s}, \check{\psi})$-system

In this section we prove that solutions of the $(\check{s}, \check{\psi})$-system (2.5) can be approximated via solutions of the phase diffusion equation (2.7).

Theorem 3.1. *Let $2 \leq m \leq n-2$. Then there exists a $C_\psi^* > 0$ such that the following is true. Let ψ^* be a solution of the phase diffusion equation* (2.7) *with*

$$\sup_{\tau' \in [0,\tau_0]} \|\psi^*(\tau)\|_{H^n_{l,u}} \leq C_\psi^*$$

and let s^* be defined by the solution $s^* = s^*(\psi^*)$ of (2.6)*. Then there exist $\delta_0 > 0$ and $C_2 > 0$ such that for all $\delta \in (0, \delta_0)$ we have solutions $(\check{s}, \check{\psi})$ of the $(\check{s}, \check{\psi})$-system* (2.5) *such that*

$$\sup_{\tau \in [0,\tau_0]} \|(\check{s}, \check{\psi})(\tau) - (s^*(\psi^*(\tau)), \psi^*(\tau))\|_{H^m_{l,u} \times H^m_{l,u}} \leq C_2 \delta^2. \tag{3.1}$$

P r o o f . Let ψ^* be a solution of the phase-diffusion equation (2.7) and let $s^* = s^*(\psi^*)$ be defined by (2.6) as above. We write a solution as approximation plus some error, i.e.

$$\check{\psi} = \psi^* + \delta^2 R_\psi \quad , \quad \check{s} = s^* + \delta^2 R_s \ .$$

We obtain

$$\begin{aligned}
\partial_\tau R_\psi &= \partial_\xi^2 R_\psi + L_\psi(R_\psi, R_s) + \delta^2 N_\psi(R_\psi, R_s) \ , \\
\partial_\tau R_s &= \partial_\xi^2 R_s - 2\delta^{-2} R_s + \delta^{-2} L_s(R_\psi, R_s) + N_s(R_\psi, R_s) + \delta^{-2} \mathrm{Res}_s \ ,
\end{aligned}$$

where

$$\begin{aligned}
L_\psi(R_\psi, R_s) &= -2\partial_\xi \left(\frac{(\partial_\xi s^*) R_\psi}{1+s^*} + \frac{(\partial_\xi R_s)\psi^*}{1+s^*} - \frac{(\partial_\xi s^*)\psi^*}{(1+s^*)^2} R_s \right) \ , \\
\delta^2 N_\psi(R_\psi, R_s) &= -2\partial_\xi \left(\frac{(\partial_\xi (s^* + \delta^2 R_s))(\psi^* + \delta^2 R_\psi)}{1 + (s^* + \delta^2 R_s)} \right) - L_\psi(R_\psi, R_s) \ , \\
L_s(R_\psi, R_s) &= -2\psi^* R_\psi - 2\psi^* s^* R_\psi - (\psi^*)^2 R_s
\end{aligned}$$

$$\begin{aligned} & -6s^*R_s - 3(s^*)^2R_s\ , \\ N_s(R_\psi, R_s) \quad = \quad & -R_\psi^2 - R_\psi^2 s^* - 2R_\psi R_s \psi^* - \delta^2 R_\psi^2 R_s \\ & -3R_s^2 - 3s^*R_s^2 - \delta^2 R_s^3\ , \\ \mathrm{Res}_s \quad = \quad & -\partial_\tau s^* + \partial_\xi^2 s^*\ . \end{aligned}$$

To be more precise, let

$$\psi^* \in C([0,\tau_0], H^n_{l,u})$$

be a solution of the phase-diffusion equation (2.7) for a $\tau_0 > 0$. Since $\partial_\tau \psi^*$ can be estimated by the right hand side of the phase-diffusion equation we have

$$\partial_\tau \psi^* \in C([0,\tau_0], H^{n-2}_{l,u})\ .$$

From the definition of s^*, we have

$$s^* \in C([0,\tau_0], H^n_{l,u})$$

and $\partial_\tau s^*, \partial_\xi^2 s^* \in C([0,\tau_0], H^{n-2}_{l,u})$. Hence, for the residual term Res_s there exists a constant C_{Res} independent of $\delta \in (0,1]$ such that

$$\sup_{\tau \in [0,\tau_0]} (\|\mathrm{Res}_s(\tau)\|_{H^{n-2}_{l,u}}) \leq C_{\mathrm{Res}}\ . \tag{3.2}$$

A natural choice to solve the equations for the error (R_ψ, R_s) would be the space $H^m_{l,u} \times H^{m+1}_{l,u}$. This choice would allow us to solve the equations as a semilinear system. However, it turns out that we have to build our proof similar to the derivation of the phase-diffusion equation (2.7) in order to gain enough powers of δ. Due to the formula for $s^* = s^*(\psi^*)$ we will consider $(R_\psi, R_s) \in H^m_{l,u} \times H^m_{l,u}$ which leads to a quasilinear system. The semilinear approach fails due to the second estimate in Lemma 3.4, for the damped part, which is sharp with respect to powers of δ.

Therefore, let us fix m for the rest of this section satisfying $2 \leq m \leq n-2$ and assume $\tau \in [0,\tau_0]$.

Lemma 3.2. *There exists a $C_1 > 0$ such that for all R_ψ and R_s and all $\delta \in (0,1]$:*

$$\begin{aligned} \|L_\psi(R_\psi, R_s)\|_{H^{m-2}_{l,u}} &\leq C_1(\|R_\psi\|_{H^m_{l,u}} + \|R_s\|_{H^m_{l,u}})\ , \\ \|L_s(R_\psi, R_s)\|_{H^m_{l,u}} &\leq C_1(\|R_\psi\|_{H^m_{l,u}} + \|R_s\|_{H^m_{l,u}})\ . \end{aligned}$$

The constant $C_1 > 0$ can be chosen to satisfy $C_1 \to 0$ if $\sup_{\tau \in [0,\tau_0]} \|\psi^*(\tau)\|_{H^n_{l,u}} \to 0$.

For all $M > 0$ there exist $\delta_1, C > 0$ such that for all $\delta \in (0,\delta_1)$ and R_ψ and R_s with $\|R_\psi\|_{H^m_{l,u}} + \|R_s\|_{H^m_{l,u}} \leq M$ we have

$$\begin{aligned} \|\delta^2 N_\psi(R_\psi, R_s)\|_{H^{m-2}_{l,u}} &\leq C\delta^2(\|R_\psi\|^2_{H^m_{l,u}} + \|R_s\|^2_{H^m_{l,u}})\ , \\ \|N_s(R_\psi, R_s)\|_{H^m_{l,u}} &\leq C(\|R_\psi\|^2_{H^m_{l,u}} + \|R_s\|^2_{H^m_{l,u}})\ . \end{aligned}$$

Proof . Follows by elementary estimates as $\|uv\|_{H^m_{l,u}} \leq \|u\|_{H^m_{l,u}} \|v\|_{H^m_{l,u}}$ for $m \geq 1$. □

Lemma 3.3. *The solution $R(\tau) = e^{(\partial_\xi^2 - a)\tau} R_0$ of*

$$\partial_\tau R = \partial_\xi^2 R - aR \quad , \quad R|_{\tau=0} = R_0$$

with $a \in \mathbb{R}$ satisfies for $r \geq 0$,

$$\|R(\tau)\|_{H_{l,u}^{m+r}} \leq Ce^{-a\tau}(1+\tau)^{-\frac{r}{2}} \|R_0\|_{H_{l,u}^m} .$$

Proof . ∂_ξ^2 is a sectorial operator in $H_{l,u}^m$ generating the analytic semigroup $e^{\partial_\xi^2 \tau}$. The estimates are standard in the theory of analytic semigroups (cf. [He81]) using multiplier theory in $H_{l,u}^m$-spaces as in [Sch99]. □

Since the nonlinear terms contain as many derivatives as the linear ones we need an optimal regularity result. We choose functions which are Hölder-continuous in time with values in $H_{l,u}^m$. We use the notation $\mu_1 = \partial_\xi^2$, $\mu_2 = \partial_\xi^2 - 2\delta^{-1}$, $c_1 = R_\psi$, and $c_2 = R_s$ and consider for $j = 1, 2$

$$\partial_\tau c_j = \mu_j c_j + g_j \tag{3.3}$$

with

$$g_j \in C^\theta([0, \tau_0], H_{l,u}^m)$$

and define

$$C_{m,j}(\tau) = \|(c_j(\tau'))_{\tau' \in [0,\tau]}\|_{C^\theta([0,\tau], H_{l,u}^m)}$$

and similarly $C_{m,g_j}(\tau)$ for a $\theta \in (0, 1)$.

Lemma 3.4. *For $r \in [0, 2]$ consider (3.3), where we additionally assume $g_j(0) = 0$ for $2\theta + r \geq 2$. Then the solutions c_j of (3.3) satisfy*

$$C_{m+r,1}(\tau) \leq CC_{m,g_1}(\tau) \quad \text{and} \quad C_{m+r,2}(\tau) \leq C\delta^{2-r} C_{m,g_2}(\tau).$$

Proof. The proof follows by direct calculation based on a classical optimal regularity result (cf. [Am95]) using the estimates on the linear semigroup from above. In detail we have: For notational simplicity we first assume $g_j(0) = 0$ and estimate with Lemma 3.3

$$\begin{aligned}
& \|\int_0^\tau e^{\mu_1(k)(\tau-\tau')} g(\tau') d\tau'\|_{H_{l,u}^{m+r}} \\
\leq \; & \|\int_0^\tau e^{\mu_1(k)(\tau-\tau')}(g(\tau') - g(\tau)) d\tau'\|_{H_{l,u}^{m+r}} + \|\int_0^\tau e^{\mu_1(k)(\tau-\tau')} d\tau' \; g(\tau)\|_{H_{l,u}^{m+r}} \\
\leq \; & C \int_0^\tau \tau'^{-r/2} \tau'^\theta d\tau' \; C_{m,g_1}(\tau) + \|\frac{1 - e^{\mu_1(k)\tau}}{\mu_1(k)} g_1(\tau)\|_{H_{l,u}^{m+r}} \\
= \; & s_1 + s_2.
\end{aligned}$$

The first term s_1 is estimated by

$$s_1 \leq C\tau^\theta \; C_{m,g_1}(\tau).$$

In order to estimate s_2 we proceed as follows

$$s_2 \leq \| \frac{1-e^{\mu_1(k)\tau}}{\mu_1(k)} \|_{H^m_{l,u} \to H^{m+r}_{l,u}} \; \|g_1(\tau)\|_{H^m_{l,u}}.$$

Since $g_1(0) = 0$ we have

$$\|g_1(\tau)\|_{H^m_{l,u}} \leq \tau^\theta C_{m,g_1}(\tau).$$

From the multiplier lemma in $H^m_{l,u}$-spaces (cf. [Sch99]) we have

$$\| \frac{1-e^{\mu_1(k)\tau}}{\mu_1(k)} \|_{H^m_{l,u} \to H^{m+r}_{l,u}} \leq \| k \mapsto \frac{1-e^{\mu_1(k)\tau}}{\mu_1(k)} (1+|k|^r) \|_{C^2_b}$$

which is estimated by

$$\begin{aligned} & \sup_{k \in \mathbb{R}} | \frac{1-e^{\mu_1(k)\tau}}{\mu_1(k)} (1+|k|^r)| \\ \leq \; & C(1+\sup_{|k|\geq 1} | \frac{|k|^r}{|k|^2}) \leq C \end{aligned}$$

and similarly for the derivatives. In a similar way we obtain

$$\begin{aligned} & \| \int_0^\tau e^{\mu_2(k)(\tau-\tau')} g(\tau') d\tau' \|_{H^{m+r}_{l,u}} \\ \leq \; & C \int_0^\tau e^{-2\delta^{-2}\tau'} \tau'^{-r/2} \tau'^\theta d\tau' \; C_{m,g_2}(\tau) + \| \frac{1-e^{\mu_2(k)\tau}}{\mu_2(k)} g_2(\tau) \|_{H^{m+r}_{l,u}} \\ \leq \; & C\delta^{2-r} \tau^\theta \; C_{m,g_1}(\tau), \end{aligned}$$

where we introduced $\delta^{-2}\tau' = \tilde{\tau}$ and estimated the integral by

$$C \int_0^\infty e^{-\tilde{\tau}} (\delta^2 \tilde{\tau})^{-r/2+\theta} \delta^2 d\tilde{\tau}.$$

For the second term we use

$$\sup_{|k|\geq 1} | \frac{k^r}{\delta^{-2}+k^2} | \leq \sup_{|k|\geq 1} | \frac{1}{(\delta^{-2}+k^2)^{1-r/2}} | \; \sup_{|k|\geq 1} | \frac{k^r}{(\delta^{-2}+k^2)^{r/2}} | \leq C\delta^{2-r}.$$

These two estimates additionally show the Hölder-continuity with exponent θ for $\tau = 0$. In a very similar fashion the Hölder-continuity for $\tau > 0$ is obtained with the same estimates in terms of δ (cf. [Am95]). The estimates for $g_j(0) \neq 0$ and $2\theta + r < 2$ follow from classical theory for semilinear equations [He81]. $\square$

Since the residual is arbitrarily smooth compared to R_ψ and R_s, it can be handled with the semilinear estimate, i.e. $r = 0$, and since $R_\psi|_{\tau=0} = R_s|_{\tau=0} = 0$ implies $L_\psi|_{\tau=0} = L_s|_{\tau=0} = N_\psi|_{\tau=0} = N_s|_{\tau=0} = 0$, Lemma 3.4 is applicable with $g_1 \in \{L_\psi, N_\psi, \mathrm{Res}_\psi\}$ and $g_2 \in \{L_s, N_s, \mathrm{Res}_s, \}$.

Hence it remains to sum up all estimates. Using (3.2), Lemma 3.2, and Lemma 3.4 gives finally the estimate

$$C_{m,1} \leq C_4(C_{m,1}+C_{m,2}) + C_5\delta^2(C_{m,1}+C_{m,2})^2 + C_{\mathrm{Res}}, \tag{3.4}$$

$$C_{m,2} \leq C_4(C_{m,1}+C_{m,2}) + C_5\delta^2(C_{m,1}+C_{m,2})^2 + C_{\mathrm{Res}} \tag{3.5}$$

with C_4 a constant satisfying $C_4 \to 0$ for $\sup_{\tau\in[0,\tau_0]} \|\psi^*(\tau)\|_{H^n_{l,u}} \to 0$, with C_5 another constant, both independent of δ. From (3.4) and (3.5) for $\sup_{\tau\in[0,\tau_0]} \|\psi^*(\tau)\|_{H^n_{l,u}} > 0$ and $\delta_0 > 0$ both sufficiently small we have

$$C_{m,1} \leq 2C_{\mathrm{Res}} \qquad \text{and} \qquad C_{m,2} \leq 2C_{\mathrm{Res}}$$

for all $\delta \in (0, \delta_0)$. Therefore, we are done. □

Remark 3.5. An alternative approach [Me98, Me99] is to consider spaces of functions $s(X,T)$, $\psi(X,T)$ that lie in a space $\mathcal{X}$ of Fourier transforms of Borel measures with bounded total variation norm. The space $\mathcal{X}$ is a Banach space with the total variation norm. We briefly describe the results of this approach, referring to [Me98, Me99] for details.

The starting point is the (s,ψ) system (2.4) or more generally the system (2.9) obtained by including higher order terms in the real Ginzburg-Landau equation. It can be shown that locally there is a one-to-one correspondence between "essential solutions" for (2.9) and essential solutions for a pseudodifferential (in time and space) equation of the form $\partial_T\psi = \partial_X^2 h + \text{h.o.t}$ where h is the function that appears in (2.7) and h.o.t are higher order terms in the same sense as in the rest of this paper — of order $O(\delta^3)$ on writing $\tau = \delta^2 T$, $\xi = \delta X$.

4. The approximation theorem for the real Ginzburg-Landau equation

It is the purpose of this final section to transfer the approximation result of Theorem 3.1, i.e. that the $(\check{s},\check{\psi})$-system (2.5) can be approximated via solutions of the phase diffusion equation (2.7) back to the real Ginzburg-Landau equation (2.1). It turns out that we cannot expect validity uniformly for all $X \in \mathbb{R}$, but validity only uniformly for all $X \in I_\delta$ with I_δ arbitrary intervals of length $O(\delta^{-r})$ for all $r < 2$.

Our starting point is the relation

$$A(X,T) = (1 + \check{s}(\delta X, \delta^2 T)) \exp\Big(i \int_0^X \check{\psi}(\delta X', \delta^2 T)\, dX'\Big)$$

which defines the solution A of the real Ginzburg-Landau equation (2.1) in terms of solutions $(\check{s},\check{\psi})$ of (2.5). These solutions are approximated by

$$A_{\mathrm{app}}(X,T) = (1 + s^*(\delta X, \delta^2 T)) \exp\Big(i \int_0^X \psi^*(\delta X', \delta^2 T)\, dX'\Big).$$

Then we obtain

$$
\begin{aligned}
& |A(X,T) - A_{\mathrm{app}}(X,T)| \\
\leq\ & \Big|(1+\check{s}(\delta X,\delta^2 T))\exp\Big(i\int_0^X \check{\psi}(\delta X',\delta^2 T)\,dX'\Big) \\
& \qquad -(1+s^*(\delta X,\delta^2 T))\exp\Big(i\int_0^X \psi^*(\delta X',\delta^2 T)\,dX'\Big)\Big| \\
\leq\ & \Big|(1+\check{s}(\delta X,\delta^2 T))\exp\Big(i\int_0^X \check{\psi}(\delta X',\delta^2 T)\,dX'\Big) \\
& \qquad -(1+\check{s}(\delta X,\delta^2 T))\exp\Big(i\int_0^X \psi^*(\delta X',\delta^2 T)\,dX'\Big)\Big| \\
& +\Big|(1+\check{s}(\delta X,\delta^2 T))\exp\Big(i\int_0^X \psi^*(\delta X',\delta^2 T)\,dX'\Big) \\
& \qquad -(1+s^*(\delta X,\delta^2 T))\exp\Big(i\int_0^X \psi^*(\delta X',\delta^2 T)\,dX'\Big)\Big| \\
\leq\ & |1+\check{s}(\delta X,\delta^2 T))|\ \Big|\exp\Big(i\int_0^X \check{\psi}(\delta X',\delta^2 T)\,dX'\Big) - \exp\Big(i\int_0^X \psi^*(\delta X',\delta^2 T)\,dX'\Big)\Big| \\
& \quad +|s^*(\delta X,\delta^2 T)) - \check{s}(\delta X,\delta^2 T))| \\
\leq\ & C\Big|\int_0^X |\check{\psi}(\delta X',\delta^2 T) - \psi^*(\delta X',\delta^2 T)|\,dX'\Big| + C\delta^2 \\
\leq\ & \Big|\int_0^X C\delta^2 dX'\Big| + C\delta^2 \leq C\delta^2(1+|X|)
\end{aligned}
$$

using the approximation result of Theorem 3.1. Thus, we have proved

Theorem 4.1. *Under the assumptions of Theorem* 3.1, *there exists a solution* $A = A(X,T)$ *of the real Ginzburg-Landau equation* (2.1) *such that the following is true. For all* $L > 0$ *and* $r \in [0,2)$ *there exist* $\delta_0 > 0$ *and* $C_2 > 0$ *such that for all* $\delta \in (0,\delta_0)$ *we have*

$$
\begin{aligned}
\sup_{T\in[0,\tau_0/\delta^2]}\ \sup_{|X|\leq L\delta^{-r}} & \Big|A(X,T) - (1+s^*(\psi^*)(\delta X,\delta^2 T)) \\
& \times\exp\Big(i\int_0^X \psi^*(\delta X',\delta^2 T)\,dX'\Big)\Big| \leq C_2\delta^{2-r}.
\end{aligned}
$$

Hence the approximation result holds uniformly on intervals larger than the natural spatial scale ($r = 1$) of the phase diffusion equation. Due to the translation invariance of the original system this holds for all intervals of length $O(\delta^{-r})$ using redefined approximations.

By taking $\check{\psi} = \delta^2$ and $\psi^* = 0$ we have to compare $e^{i\delta^2 X}$ with 1 which shows that estimates uniformly valid for all $X \in \mathbb{R}$ cannot be expected. A uniform estimate can only be expected for $\check{\psi}$ and ψ^* spatially localized, for instance in spaces like $H^m(n)$ with $n \geq 2$, equipped with $\|u\|_{H^m(n)} = \|u\rho^n\|_{H^m}$ with $\rho(x) = \sqrt{1+x^2}$. We claim the validity of the approximation result for $(\check{\psi}, \check{s}) \in H^m(2) \times H^m(2)$, too.

Remark 4.2. By adding higher order terms in (ψ^*, s^*), the approximation in Theorem 4.1 can be improved to hold uniformly on space intervals with length of order $O(\delta^{-p})$ with $p \geq 0$ arbitrary. Indeed, it suffices to include all terms in (ψ^*, s^*) of order $O(\delta^{\tilde{p}})$ for some fixed $\tilde{p} > p$. In order to do so we make the ansatz

$$\begin{aligned} \check{\psi} &= \psi^* + \delta\psi_1^* + \delta^2\psi_2^* + \ldots + \delta^{\tilde{p}}\psi_{\tilde{p}}^*, \\ \check{s} &= s^* + \delta s_1^* + \delta^2 s_2^* + \ldots + \delta^{\tilde{p}}\psi_{\tilde{p}}^*. \end{aligned}$$

The governing equations for the ψ_j^* and s_j^* are obtained by equating the coefficients of δ^{j+2} and δ^j to zero. They are linear with respect to ψ_j^* and s_j^* for fixed j. For instance, we obtain for s_1^* the equation

$$0 = -2s_1^* - 2\psi_1^*\psi^* - 2\psi_1^*\psi^* s^* - (\psi^*)^2 s_1^* - 6s_1^* s^* - 3(s^*)^2 s_1^*.$$

Hence the solutions of these equations exist as long as the solutions of the phase diffusion equation exist, again under the assumption that ψ^* is sufficiently small.

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